

# **Geodesics in standard stationary spacetimes and Lagrangian systems**

Anna Valeria Germinario  
Università di Bari – Italy

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**Standard stationary spacetime:**  $(L, \langle \cdot, \cdot \rangle_L)$  Lorentzian manifold,  $L = M \times \mathbb{R}$  and

$$\langle \cdot, \cdot \rangle_L = \langle \cdot, \cdot \rangle + 2\langle \delta(x), \cdot \rangle dt - \beta(x) dt^2,$$

$(M, \langle \cdot, \cdot \rangle)$  finite dimensional, connected Riemannian manifold,  $\delta$  smooth vector field on  $M$  and  $\beta : M \rightarrow \mathbb{R}$ , smooth strictly positive function.

If  $\delta \equiv 0$ , then  $L$  is called **standard static**.

Hypothesis:

- $(M, \langle \cdot, \cdot \rangle)$  complete, at least  $C^3$  Riemannian manifold.

**Conservation laws:** if a curve  $z = (x, t) : I \rightarrow L$ ,  $I \subset \mathbb{R}$  interval, is a geodesic, then  $E \in \mathbb{R}$  and  $K \in \mathbb{R}$  exist such that

$$E = \frac{1}{2} \langle \dot{z}, \dot{z} \rangle_L \quad K = -\langle \partial t, \dot{z} \rangle_L = \beta(x)\dot{t} - \langle \delta(x), \dot{x} \rangle.$$

Variational approach for Lorentzian geodesics:

every geodesic  $z : [a, b] \rightarrow L$  is a critical point of the (strongly indefinite) functional

$$f(z) = \int_a^b \langle \dot{z}, \dot{z} \rangle_L ds$$

defined on a manifold of curves.

- V. Benci, D. Fortunato, F. Giannoni (Ann. Inst. H. Poincaré Anal. Non Linéaire, 1991) and others: reduction to the study of a Riemannian functional.

- M. Sánchez, (Nonlinear Anal., 1999) in the [standard static case](#): relation between geodesics and Lagrangian systems.

**Theorem.** Let  $(L, \langle \cdot, \cdot \rangle_L)$  be a st. static sp. A curve  $z = (x, t) : I \rightarrow L$ ,  $I \subset \mathbb{R}$  interval, is a geodesic such that  $\beta(x)\dot{t} = K$ , if and only if  $x$  solves

$$D_s \dot{x} + \frac{1}{2} K^2 \nabla V(x) = 0 \quad V = -\frac{1}{\beta}$$

( $D_s$  covariant derivative and  $\nabla$  gradient with respect to  $\langle \cdot, \cdot \rangle$ ). Moreover

$$\frac{1}{2} \langle \dot{z}(s), \dot{z}(s) \rangle_L = \frac{1}{2} \langle \dot{x}(s), \dot{x}(s) \rangle + \frac{1}{2} K^2 V(x(s)) \quad \forall s \in I.$$

Results in A.V.G. (J. Differential Equations, 2007):

- a correspondence between geodesics and Lagrangian systems holds also in the standard stationary case if
  1.  $M$  is endowed by a **new Riemannian metric** (depending on the coefficients  $\delta$  and  $\beta$ );
  2. a term representing the **action of a magnetic field** is added in the previous differential equation.

New (complete) metric on  $M$ : set  $\langle \cdot, \cdot \rangle_1$  by

$$\langle u, v \rangle_1 = \langle u, v \rangle + \frac{1}{\beta(x)} \langle \delta(x), u \rangle \langle \delta(x), v \rangle \quad x \in M, u, v \in T_x M.$$

For any  $v \in T_x M$ , it is

$$\langle v, v \rangle_1 = \langle (I + P(x))[v], v \rangle \quad P(x)[v] = \frac{1}{\beta(x)} \langle \delta(x), v \rangle \delta(x),$$

$I$  identity,  $P(x) : T_x M \rightarrow T_x M$  linear, self-adjoint, positive operator.

For any  $x \in M$ , set

$$V(x) = -\frac{1}{\beta(x)} \quad A(x) = \frac{\delta(x)}{\beta(x) + \langle \delta(x), \delta(x) \rangle}.$$

Let  $F^1$  be the curl of  $A$ :

$$F^1(X, Y) = \langle D_X^1 A, Y \rangle_1 - \langle X, D_Y^1 A \rangle_1$$

( $X, Y$  vector fields on  $M$ ,  $D^1$  Levi-Civita connection of  $(M, \langle \cdot, \cdot \rangle_1)$ ).

New equation:

$$D_s^1 \dot{x} + \frac{1}{2} K^2 \nabla^1 V(x) = K \hat{F}^1(x)[\dot{x}] \quad (1)$$

$\nabla^1$  gradient with resp. to  $\langle \cdot, \cdot \rangle_1$ ,  $\hat{F}^1 : TM \rightarrow TM$  linear map  $\langle \cdot, \cdot \rangle_1$ -associated to  $F^1$  ( $F^1(x)[u, v] = \langle \hat{F}^1(x)[u], v \rangle_1$ ,  $x \in M$ ,  $u, v \in T_x M$ ).

**Theorem.** Let  $(L, \langle \cdot, \cdot \rangle_L)$  be a st. stationary sp.,  $I \subset \mathbb{R}$  an interval.

If  $z = (x, t) : I \rightarrow L$  is a geodesic such that  $\beta(x)\dot{t} - \langle \delta(x), \dot{x} \rangle = K$  for some  $K \in \mathbb{R}$ , then  $x : I \rightarrow \mathbb{R}$  solves (1).

Every geodesic  $z = (x, t) : I \rightarrow L$  for  $\langle \cdot, \cdot \rangle_L$  can be obtained by a sol.  $x : I \rightarrow M$  of (1) for some  $K \in \mathbb{R}$ ,  $t$  by  $\dot{t} = (K + \langle \delta(x), \dot{x} \rangle) / \beta(x)$ .

Moreover

$$\frac{1}{2} \langle \dot{z}, \dot{z} \rangle_L = \frac{1}{2} \langle \dot{x}, \dot{x} \rangle_1 + \frac{1}{2} K^2 V(x) \quad \text{on } I.$$

The proof follows from the [computation of the geodesic equations](#) on  $L$ .

A geodesic  $z = (x, t) : [a, b] \rightarrow L$  is a critical point of

$$\begin{aligned} f(z) &= \int_a^b \langle \dot{z}, \dot{z} \rangle_L ds \\ &= \int_a^b \left( \langle \dot{x}, \dot{x} \rangle_1 - \beta(x) (t - \langle A(x), \dot{x} \rangle_1)^2 \right) ds. \end{aligned}$$

Differentiating  $f$  gives

$$\begin{cases} \frac{d}{ds} (\beta(x) (t - \langle A(x), \dot{x} \rangle_1)) = 0 \\ D_s^1 \dot{x} + \frac{1}{2} (t - \langle A(x), \dot{x} \rangle_1)^2 \nabla^1 \beta(x) = \beta(x) (t - \langle A(x), \dot{x} \rangle_1) \hat{F}^1(x) [\dot{x}]. \end{cases}$$



**Applications:** geodesics with fixed  $E$  (in a suitable real interval) and  $K = \sqrt{2}$  (as in the static case, normalization of the coefficient of  $\nabla^1 V$ ).

**Boundary conditions:** geod. from a point to a line or periodic trajectories.

- sol.  $x : [0, a] \rightarrow M$  of

$$D_s^1 \dot{x} + \nabla^1 V(x) = \sqrt{2} \hat{F}^1(x)[\dot{x}]$$

joining two fixed points  $x_0, x_1 \in M$  correspond to geod.  $z = (x, t) : [0, a] \rightarrow L$  joining a point  $z_0 = (x_0, t_0)$  to a line  $(x_1, s) \subset L$ , where

$$t(s) = t_0 + \int_0^s \frac{\sqrt{2} + \langle \delta(x), \dot{x} \rangle}{\beta(x)} d\tau.$$

Case studied by R. Bartolo, D. Fortunato, F. Giannoni, A. Masiello, P. Piccione, M. Sánchez and others.

**Theorem.** Let  $(L, \langle \cdot, \cdot \rangle_L)$  be a st. stationary sp. If

- $\bar{A} \in \mathbb{R}$  exists such that

$$\sup_{x \in M} |A(x)|_1 = \sup_{x \in M} \frac{|\delta(x)|}{\sqrt{\beta(x)(\beta(x) + \langle \delta(x), \delta(x) \rangle)}} = \bar{A}.$$

Then, for any  $E \in \mathbb{R}$  with

$$E > \bar{\beta} + \bar{A}^2 \quad \text{where} \quad \bar{\beta} = \sup_{x \in M} \left( -\frac{1}{\beta(x)} \right)$$

and for any  $x_0, x_1 \in M$ ,  $x_0 \neq x_1$ ,  $t_0 \in \mathbb{R}$ , a geodesic  $z = (y, t) : [0, a] \rightarrow L$  exists joining the point  $(x_0, t_0) \in L$  to the line  $(x_1, s) \subset L$ , such that

$$\frac{1}{2} \langle \dot{z}, \dot{z} \rangle_L = E \quad \text{and} \quad \beta(y)t - \langle \delta(y), \dot{y} \rangle = \sqrt{2}.$$

## Multiplicity:

- if  $M$  is not contractible in itself, a **sequence**  $(z_m)$ ,  $z_m = (y_m, t_m) : [0, a_m] \rightarrow M$  of geodesics exists. Their **arrival times**  $t_m(a_m)$  verify

$$\lim_{m \rightarrow +\infty} t_m(a_m) = +\infty$$

when  $\beta$  is **bounded from above** and, denoted by

$$\bar{A}_1 = \sup_{x \in M} \frac{|\delta(x)|}{\sqrt{\beta(x) + \langle \delta(x), \delta(x) \rangle}} \quad N = \sup_{x \in M} \beta(x),$$

when

- $\bar{A}_1 < \sqrt{EN + 1}$  for any possible  $E \leq 0$ ;
- $\bar{A}_1 < 1/\sqrt{EN + 1}$  for any  $E > 0$ .

The **boundedness condition** is satisfied when  $\beta \geq \nu$ , for some  $\nu > 0$  or  $|\delta|/\beta$  is bounded.

$E > \bar{\beta} + \bar{A}^2$ : If  $\beta$  is **bounded from above** then  $\bar{\beta} < 0$  (thus in some cases negative  $E$  are allowed). If  $\bar{\beta} \geq 0$ , unbounded interval of strictly positive  $E$ .

The theorem contains, as **particular cases**, some of the th. in

- D. Fortunato, F. Giannoni, A. Masiello, J. Geom. Phys., 1995 (case  $E = 0$  on st. stationary sp.);
- R. Bartolo, A.V.G., M. Sánchez, Differential Geom. Appl., 2002, (st. static sp.,  $E > \bar{\beta}$ ).

## Periodic trajectories:

Periodic trajectory of universal period  $T$  and proper period  $a > 0$ : a geodesic  $z = (x, t) : [0, a] \rightarrow L$  such that

$$x(a) = x(0) \quad \dot{x}(a) = \dot{x}(0) \quad t(a) = t(0) + T \quad \dot{t}(a) = \dot{t}(0).$$

- periodic sol. (i.e.  $x : [0, a] \rightarrow M$  such that  $x(0) = x(a)$ ,  $\dot{x}(0) = \dot{x}(a)$ ) of

$$D_s^1 \dot{x} + \nabla^1 V(x) = \sqrt{2} \hat{F}^1(x) [\dot{x}]$$

give rise to periodic trajectories  $z = (x, t) : [0, a] \rightarrow L$ .

Set  $t(0) = 0$ , thus

$$T = t(a) = \int_0^a \frac{\sqrt{2} + \langle \delta(x), \dot{x} \rangle}{\beta(x)} ds.$$

**Geometrically distinct** periodic trajectories: if they have **different ranges**. Taking  $t(0) = 0$  avoids obtaining trajectories having the same spatial components and with temporal components differing by a constant.

**Theorem.** Let  $(L, \langle \cdot, \cdot \rangle_L)$  be a st. stationary sp. with **compact**  $M$ .  
If

- $M$  is not contractible in itself and its fundamental group  $\pi_1(M)$  is finite or it has infinitely many conjugacy classes.

Then, for any  $E > \bar{\beta} + \bar{A}^2$ , one non-trivial  $t$ -periodic trajectory  $z = (y, t) : [0, a] \rightarrow L$  exists, such that

$$\frac{1}{2} \langle \dot{z}, \dot{z} \rangle_L = E \quad \text{and} \quad \beta(y)t - \langle \delta(y), \dot{y} \rangle = \sqrt{2}.$$

Fixing  $E$  and the value  $K = \sqrt{2}$  gives a **multiplicity result**: periodic traj. for different values of  $E$  are geometrically distinct.

$M$  **non-compact**: further hyp. at infinity (existence of a convex at infinity function as in V. Benci, D. Fortunato, Proc. of “Variational Methods”, 1988).

The theorem **extends** some results in

- A. Candela, Ann. Mat. Pura Appl., 1996, ( $E = 0$ , st. stationary sp.).

In the non-compact case, **weaker assumptions** than

- A. Masiello, Nonlinear Anal., 1992 (st. stationary sp., traj. with fixed  $T$ ).

Proof: find fixed energy solutions of

$$D_s^1 \dot{x} + \nabla^1 V(x) = \sqrt{2} \hat{F}^1(x)[\dot{x}].$$

Let  $(M, \langle \cdot, \cdot \rangle)$  be a complete Riem. manifold,  $V : M \rightarrow \mathbb{R}$  a smooth function,  $B$  smooth vector field on  $M$ ,  $F$  curl of  $B$ .

$$D_s \dot{x} + \nabla V(x) = \hat{F}(x)[\dot{x}] \quad (2)$$

where  $\hat{F} : TM \rightarrow TM$  is the linear map  $\langle \cdot, \cdot \rangle$ -associated to  $F$ .

Fixed  $E_x \in \mathbb{R}$ , find solutions  $x : I \rightarrow M$ ,  $I \subset \mathbb{R}$  interval, of (2) s.t.

$$E = \frac{1}{2} \langle \dot{x}(s), \dot{x}(s) \rangle + V(x(s)) \quad \forall s \in I.$$



If  $F \equiv 0$ , **classical principle**: solutions of (2) with fixed energy  $E$ , are, up to reparametrizations, geodesics with respect to a **Jacobi metric**

$$\langle \cdot, \cdot \rangle_E = (E - V(x)) \langle \cdot, \cdot \rangle$$

(Riemannian in a neighbourhood of  $x \in M$ , if  $E > V(x)$ ).

If  $F \neq 0$ ,

- $V \in C^1(M, \mathbb{R})$  is **bounded from above**;
- $B$  is **bounded**.

$$\bar{V} = \sup_{x \in M} V(x) \quad \bar{B} = \sup_{x \in M} |B(x)|$$

**Proposition.** Let  $E > \bar{V}$ . If  $x \in C^2([0, 1], M)$  is a non-constant solution of

$$(E - V(x))D_s^E \dot{x} = \sqrt{\frac{1}{2}\langle \dot{x}, \dot{x} \rangle_E} \hat{F}(x)[\dot{x}] \quad (3)$$

( $D_s^E$  covariant derivative with respect to  $\langle \cdot, \cdot \rangle_E$ ), then  $a > 0$  and a reparametrization  $y \in C^2([0, a], M)$  of  $x$  exist, solving

$$D_s \dot{x} + \nabla V(x) = \hat{F}(x)[\dot{x}]$$

and having energy  $E$ .

Eq. (3) is **invariant by affine rep.**  $a_1 s + a_2$  if  $a_1 \geq 0$ .

Eq. (3) has a **variational structure**: its associated functional is

$$G(x) = \sqrt{2 \int_0^1 \langle \dot{x}, \dot{x} \rangle_E ds} + \int_0^1 \langle B(x), \dot{x} \rangle ds$$

$x$  varying in a **suitable manifold of curves**.

For some fixed  $x_0, x_1$  in  $M$ ,  $x_0 \neq x_1$ , consider

$$\Omega^1(x_0, x_1, M) = \{x \in H^1([0, 1], M) \mid x(0) = x_0, x(1) = x_1\}$$

$$\Lambda^1(M) = \{x \in H^1([0, 1], M) \mid x(0) = x(1)\}.$$

Critical points of  $G(x)$ ,  $x \in \Omega^1(x_0, x_1, M)$  exist, under the inequality

$$E > \bar{V} + \frac{\bar{B}^2}{2}.$$

( $G$  has minimum. If  $M$  is not contractible, multiplicity of critical points.)

If  $x \in \Lambda^1(M)$ , the minimum of  $G$  is 0: reinforce topological assumptions about  $M$  (thus  $\Lambda^1(M)$  has subsets of arbitrarily large Ljusternik–Schnirelman category).

Work in progress:

if  $M$  is not complete and has boundary  $\partial M$ ,

- extend these results (under a suitable notion of convexity for  $\partial M$ , with respect to the field  $A$ );
- study the convexity of  $\partial L = \partial M \times \mathbb{R}$  using metric  $\langle \cdot, \cdot \rangle_1$ .